

Universal non-perturbative features in the thermal behavior of non-Abelian gauge theories

Marco Panero

University of Turin and INFN, Turin, Italy

NIC, DESY Zeuthen
30 November 2015



Outline

Introduction

Thermodynamics in the confining phase: comparing $SU(2)$ and $SU(3)$ Yang-Mills theories

Thermodynamics in the deconfined phase: comparing $SU(N)$ and G_2 theories

Conclusions

Based on:

- ★ M. Caselle, A. Nada and M. P., JHEP **07** (2015) 143, [arXiv:1505.01106](#)
- ★ M. Bruno, M. Caselle, M. P. and R. Pellegrini, JHEP **03** (2015) 057, [arXiv:1409.8305](#)



Outline

Introduction

Thermodynamics in the confining phase: comparing $SU(2)$ and $SU(3)$ Yang-Mills theories

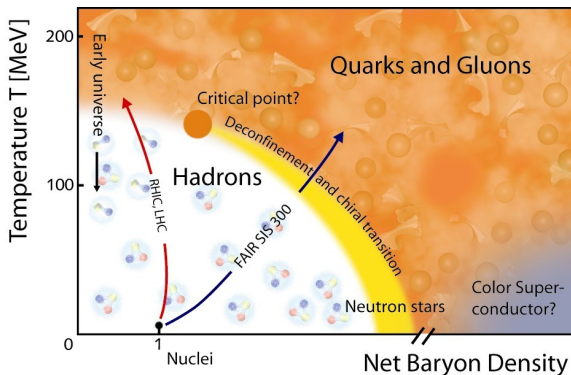
Thermodynamics in the deconfined phase: comparing $SU(N)$ and G_2 theories

Conclusions



Strong nuclear interactions under extreme conditions

Quantum Chromodynamics is believed to possess a rich phase structure



[N. Cabibbo and G. Parisi, 1975], [J. C. Collins and M. J. Perry, 1975]



Why is it important?



Adapted from [Z. Weiner, 2010]



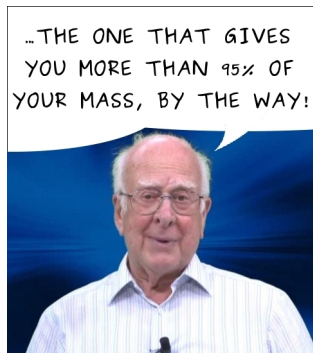
Why is it important?

- ▶ **An important test of QCD, one of the building blocks of the Standard Model**
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model



- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [**B. Müller, 2013**]
 - ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
 - ▶ Very rich physics, involving several non-trivial theoretical problems
 - ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
 - ★ Simultaneously relativistic and strongly coupled
 - ★ A liquid that cools into a gas
 - ★ “The most perfect” liquid—shear viscosity close to its quantum limit [G. Policastro, D. T. Son and A. O. Starinets, 2001]
 - ★ Fast thermalization—close to limit set by causality
 - ★ It creates its own vacuum state to exist
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
 - ★ Simultaneously relativistic and strongly coupled
 - ★ A liquid that cools into a gas
 - ★ “The most perfect” liquid—shear viscosity close to its quantum limit [G. Policastro, D. T. Son and A. O. Starinets, 2001]
 - ★ Fast thermalization—close to limit set by causality
 - ★ It creates its own vacuum state to exist
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [**B. Müller, 2013**]
 - ★ Simultaneously relativistic and strongly coupled
 - ★ A liquid that cools into a gas
 - ★ “The most perfect” liquid—shear viscosity close to its quantum limit [**G. Policastro, D. T. Son and A. O. Starinets, 2001**]
 - ★ Fast thermalization—close to limit set by causality
 - ★ It creates its own vacuum state to exist
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [**B. Müller, 2013**]
 - ★ Simultaneously relativistic and strongly coupled
 - ★ A liquid that cools into a gas
 - ★ “The most perfect” liquid—shear viscosity close to its quantum limit [**G. Policastro, D. T. Son and A. O. Starinets, 2001**]
 - ★ Fast thermalization—close to limit set by causality
 - ★ It creates its own vacuum state to exist
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [**E. Shuryak, 2009**]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [**B. Müller, 2013**]
 - ★ Simultaneously relativistic and strongly coupled
 - ★ A liquid that cools into a gas
 - ★ “The most perfect” liquid—shear viscosity close to its quantum limit [**G. Policastro, D. T. Son and A. O. Starinets, 2001**]
 - ★ Fast thermalization—close to limit set by causality
 - ★ It creates its own vacuum state to exist
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [**E. Shuryak, 2009**]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems \Rightarrow Requires a combination of different tools: weak-coupling expansions, lattice simulations, effective theories, phenomenological models, holography, ...
- ▶ The focus of a *large, successful and long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Why is it important?

- ▶ An important test of QCD, one of the building blocks of the Standard Model
- ▶ Temperatures $\gtrsim 200$ MeV realized in nature until about $1 \mu\text{s}$ after the Big Bang; the early Universe cooling rate depends on the QCD equation of state (EoS)
- ▶ Cold and dense QCD matter probably exists in compact stars
- ▶ The quark-gluon plasma (QGP) has very peculiar properties [B. Müller, 2013]
- ▶ Connections to seemingly distant physical systems: superfluids, ultracold atoms, fermionic condensed matter systems, black holes, ... [E. Shuryak, 2009]
- ▶ Very rich physics, involving several non-trivial theoretical problems \Rightarrow Requires a combination of different tools: weak-coupling expansions, lattice simulations, effective theories, phenomenological models, holography, ...
- ▶ The focus of a *large, successful* and *long-lasting* experimental programme through heavy-ion collisions (BNL, LHC, GSI, JINR, ...)



Motivation to study finite-temperature QCD and QCD-like theories on the lattice

- ▶ Robust *ab initio* theoretical predictions for strong interactions under the conditions probed in heavy-ion collisions
- ▶ Weak-coupling expansions in thermal QCD become accurate only at high temperatures, fail to capture the physics of long-wavelength modes in the deconfined phase [A. Linde, 1980] and the dynamics in the hadronic phase
- ▶ The lattice regularization [K. G. Wilson, 1974] provides *the only known* mathematically well-defined, non-perturbative formulation of QCD
- ▶ Thanks to steady theoretical, algorithmic and computer-power progress, lattice computations are now producing accurate **numerical** predictions for QCD at finite temperature
- ▶ The lattice investigation of finite-temperature QCD-like theories can lead to **analytical** insight of the underlying physics



Motivation to study finite-temperature QCD and QCD-like theories on the lattice

- ▶ Robust *ab initio* theoretical predictions for strong interactions under the conditions probed in heavy-ion collisions
- ▶ Weak-coupling expansions in thermal QCD become accurate only at high temperatures, fail to capture the physics of long-wavelength modes in the deconfined phase [**A. Linde, 1980**] and the dynamics in the hadronic phase
- ▶ The lattice regularization [K. G. Wilson, 1974] provides *the only known* mathematically well-defined, non-perturbative formulation of QCD
- ▶ Thanks to steady theoretical, algorithmic and computer-power progress, lattice computations are now producing accurate **numerical** predictions for QCD at finite temperature
- ▶ The lattice investigation of finite-temperature QCD-like theories can lead to **analytical** insight of the underlying physics



Motivation to study finite-temperature QCD and QCD-like theories on the lattice

- ▶ Robust *ab initio* theoretical predictions for strong interactions under the conditions probed in heavy-ion collisions
- ▶ Weak-coupling expansions in thermal QCD become accurate only at high temperatures, fail to capture the physics of long-wavelength modes in the deconfined phase [A. Linde, 1980] and the dynamics in the hadronic phase
- ▶ The lattice regularization [K. G. Wilson, 1974] provides *the only known* mathematically well-defined, non-perturbative formulation of QCD
- ▶ Thanks to steady theoretical, algorithmic and computer-power progress, lattice computations are now producing accurate **numerical** predictions for QCD at finite temperature
- ▶ The lattice investigation of finite-temperature QCD-like theories can lead to **analytical** insight of the underlying physics



Motivation to study finite-temperature QCD and QCD-like theories on the lattice

- ▶ Robust *ab initio* theoretical predictions for strong interactions under the conditions probed in heavy-ion collisions
- ▶ Weak-coupling expansions in thermal QCD become accurate only at high temperatures, fail to capture the physics of long-wavelength modes in the deconfined phase [A. Linde, 1980] and the dynamics in the hadronic phase
- ▶ The lattice regularization [K. G. Wilson, 1974] provides *the only known* mathematically well-defined, non-perturbative formulation of QCD
- ▶ Thanks to steady theoretical, algorithmic and computer-power progress, lattice computations are now producing accurate **numerical** predictions for QCD at finite temperature
- ▶ The lattice investigation of finite-temperature QCD-like theories can lead to **analytical** insight of the underlying physics



Motivation to study finite-temperature QCD and QCD-like theories on the lattice

- ▶ Robust *ab initio* theoretical predictions for strong interactions under the conditions probed in heavy-ion collisions
- ▶ Weak-coupling expansions in thermal QCD become accurate only at high temperatures, fail to capture the physics of long-wavelength modes in the deconfined phase [A. Linde, 1980] and the dynamics in the hadronic phase
- ▶ The lattice regularization [K. G. Wilson, 1974] provides *the* only known mathematically well-defined, non-perturbative formulation of QCD
- ▶ Thanks to steady theoretical, algorithmic and computer-power progress, lattice computations are now producing accurate **numerical** predictions for QCD at finite temperature
- ▶ The lattice investigation of finite-temperature QCD-like theories can lead to **analytical** insight of the underlying physics



Outline

Introduction

Thermodynamics in the confining phase: comparing $SU(2)$ and $SU(3)$ Yang-Mills theories

Thermodynamics in the deconfined phase: comparing $SU(N)$ and G_2 theories

Conclusions



Glueball gas thermodynamics

- ▶ The physical states in the confining, low-temperature phase are color-singlet hadrons
- ▶ For non-supersymmetric, pure-gluon SU(N) theories, these hadrons are glueballs
- ▶ Glueballs are classified according to their J^{PC} quantum numbers, have a finite mass gap, and are weakly coupled
- ▶ The equation of state can be described in terms of a gas of massive, free bosons; for example, for the pressure:

$$p(T) = -\frac{\partial F}{\partial V} = \lim_{V \rightarrow \infty} \frac{T}{V} \ln Z = \sum_i (2J_i + 1) \rho(m_i; T)$$

where $\rho(m_i; T)$ denotes the contribution from each degree of freedom of a state of mass m_i :

$$\rho(m_i; T) = \frac{T^4}{2} \left(\frac{m_i}{T\pi} \right)^2 \sum_{n=1}^{\infty} \frac{1}{n^2} K_2 \left(n \frac{m_i}{T} \right)$$



Glueball gas thermodynamics

- ▶ The physical states in the confining, low-temperature phase are color-singlet hadrons
- ▶ For non-supersymmetric, pure-gluon SU(N) theories, these hadrons are glueballs
- ▶ Glueballs are classified according to their J^{PC} quantum numbers, have a finite mass gap, and are weakly coupled
- ▶ The equation of state can be described in terms of a gas of massive, free bosons; for example, for the pressure:

$$p(T) = -\frac{\partial F}{\partial V} = \lim_{V \rightarrow \infty} \frac{T}{V} \ln Z = \sum_i (2J_i + 1) \rho(m_i; T)$$

where $\rho(m_i; T)$ denotes the contribution from each degree of freedom of a state of mass m_i :

$$\rho(m_i; T) = \frac{T^4}{2} \left(\frac{m_i}{T\pi} \right)^2 \sum_{n=1}^{\infty} \frac{1}{n^2} K_2 \left(n \frac{m_i}{T} \right)$$



Glueball gas thermodynamics

- ▶ The physical states in the confining, low-temperature phase are color-singlet hadrons
- ▶ For non-supersymmetric, pure-gluon SU(N) theories, these hadrons are glueballs
- ▶ Glueballs are classified according to their J^{PC} quantum numbers, have a finite mass gap, and are weakly coupled
- ▶ The equation of state can be described in terms of a gas of massive, free bosons; for example, for the pressure:

$$p(T) = -\frac{\partial F}{\partial V} = \lim_{V \rightarrow \infty} \frac{T}{V} \ln Z = \sum_i (2J_i + 1) \rho(m_i; T)$$

where $\rho(m_i; T)$ denotes the contribution from each degree of freedom of a state of mass m_i :

$$\rho(m_i; T) = \frac{T^4}{2} \left(\frac{m_i}{T\pi} \right)^2 \sum_{n=1}^{\infty} \frac{1}{n^2} K_2 \left(n \frac{m_i}{T} \right)$$



Glueball gas thermodynamics

- ▶ The physical states in the confining, low-temperature phase are color-singlet hadrons
- ▶ For non-supersymmetric, pure-gluon SU(N) theories, these hadrons are glueballs
- ▶ Glueballs are classified according to their J^{PC} quantum numbers, have a finite mass gap, and are weakly coupled
- ▶ The equation of state can be described in terms of a gas of massive, free bosons; for example, for the pressure:

$$p(T) = -\frac{\partial F}{\partial V} = \lim_{V \rightarrow \infty} \frac{T}{V} \ln Z = \sum_i (2J_i + 1) p(m_i; T)$$

where $p(m_i; T)$ denotes the contribution from each degree of freedom of a state of mass m_i :

$$p(m_i; T) = \frac{T^4}{2} \left(\frac{m_i}{T\pi} \right)^2 \sum_{n=1}^{\infty} \frac{1}{n^2} K_2 \left(n \frac{m_i}{T} \right)$$



A string model for the spectrum

- ▶ How many glueball states should be included in the sum?
 - ★ At temperatures much lower than the deconfinement transition temperature T_c , the sum is completely dominated by the lightest glueballs
 - ★ At temperatures closer to T_c , the contribution of heavier states becomes non-negligible
- ▶ While the lattice determination of heavy states in the spectrum is challenging [M. Teper, 1998], the spectral density $\hat{\rho}(m)$ can be predicted by modelling glueballs as “rings of glue” [N. Isgur and J. E. Paton, 1984]
- ▶ The effective description of confining flux tubes as open strings [M. Lüscher, 1980] suggests to model the glueballs as closed bosonic strings, leading to:

$$\hat{\rho}(m) = \frac{1}{m} \left(\frac{2\pi T_H}{3m} \right)^3 \exp(m/T_H)$$

where T_H is the Hagedorn temperature [R. Hagedorn, 1965]

- ▶ The total contribution to the pressure, including contributions from the lightest states and from a continuous spectrum above a threshold $m_{\text{threshold}} = 2m_{0^{++}}$, reads

$$p(T) = \sum_{m_i < m_{\text{threshold}}} (2J_i + 1) p(m_i; T) + n_C \int_{m_{\text{threshold}}}^{\infty} dm' \hat{\rho}(m') p(m'; T)$$

where n_C accounts for the charge-conjugation multiplicity



A string model for the spectrum

- ▶ How many glueball states should be included in the sum?
 - ★ At temperatures much lower than the deconfinement transition temperature T_c , the sum is completely dominated by the lightest glueballs
 - ★ At temperatures closer to T_c , the contribution of heavier states becomes non-negligible
- ▶ While the lattice determination of heavy states in the spectrum is challenging [M. Teper, 1998], the spectral density $\hat{\rho}(m)$ can be predicted by modelling glueballs as “rings of glue” [N. Isgur and J. E. Paton, 1984]
- ▶ The effective description of confining flux tubes as open strings [M. Lüscher, 1980] suggests to model the glueballs as closed bosonic strings, leading to:

$$\hat{\rho}(m) = \frac{1}{m} \left(\frac{2\pi T_H}{3m} \right)^3 \exp(m/T_H)$$

where T_H is the Hagedorn temperature [R. Hagedorn, 1965]

- ▶ The total contribution to the pressure, including contributions from the lightest states and from a continuous spectrum above a threshold $m_{\text{threshold}} = 2m_{0^{++}}$, reads

$$p(T) = \sum_{m_i < m_{\text{threshold}}} (2J_i + 1) p(m_i; T) + n_C \int_{m_{\text{threshold}}}^{\infty} dm' \hat{\rho}(m') p(m'; T)$$

where n_C accounts for the charge-conjugation multiplicity



A string model for the spectrum

- ▶ How many glueball states should be included in the sum?
 - ★ At temperatures much lower than the deconfinement transition temperature T_c , the sum is completely dominated by the lightest glueballs
 - ★ At temperatures closer to T_c , the contribution of heavier states becomes non-negligible
- ▶ While the lattice determination of heavy states in the spectrum is challenging [M. Teper, 1998], the spectral density $\hat{\rho}(m)$ can be predicted by modelling glueballs as “rings of glue” [N. Isgur and J. E. Paton, 1984]
- ▶ The effective description of confining flux tubes as open strings [M. Lüscher, 1980] suggests to model the glueballs as closed bosonic strings, leading to:

$$\hat{\rho}(m) = \frac{1}{m} \left(\frac{2\pi T_H}{3m} \right)^3 \exp(m/T_H)$$

where T_H is the Hagedorn temperature [R. Hagedorn, 1965]

- ▶ The total contribution to the pressure, including contributions from the lightest states and from a continuous spectrum above a threshold $m_{\text{threshold}} = 2m_{0^{++}}$, reads

$$p(T) = \sum_{m_i < m_{\text{threshold}}} (2J_i + 1) p(m_i; T) + n_C \int_{m_{\text{threshold}}}^{\infty} dm' \hat{\rho}(m') p(m'; T)$$

where n_C accounts for the charge-conjugation multiplicity



A string model for the spectrum

- ▶ How many glueball states should be included in the sum?
 - ★ At temperatures much lower than the deconfinement transition temperature T_c , the sum is completely dominated by the lightest glueballs
 - ★ At temperatures closer to T_c , the contribution of heavier states becomes non-negligible
- ▶ While the lattice determination of heavy states in the spectrum is challenging [M. Teper, 1998], the spectral density $\hat{\rho}(m)$ can be predicted by modelling glueballs as “rings of glue” [N. Isgur and J. E. Paton, 1984]
- ▶ The effective description of confining flux tubes as open strings [M. Lüscher, 1980] suggests to model the glueballs as closed bosonic strings, leading to:

$$\hat{\rho}(m) = \frac{1}{m} \left(\frac{2\pi T_H}{3m} \right)^3 \exp(m/T_H)$$

where T_H is the Hagedorn temperature [R. Hagedorn, 1965]

- ▶ The total contribution to the pressure, including contributions from the lightest states and from a continuous spectrum above a threshold $m_{\text{threshold}} = 2m_{0^{++}}$, reads

$$p(T) = \sum_{m_i < m_{\text{threshold}}} (2J_i + 1) p(m_i; T) + n_C \int_{m_{\text{threshold}}}^{\infty} dm' \hat{\rho}(m') p(m'; T)$$

where n_C accounts for the charge-conjugation multiplicity



A string model for the spectrum

- ▶ How many glueball states should be included in the sum?
 - ★ At temperatures much lower than the deconfinement transition temperature T_c , the sum is completely dominated by the lightest glueballs
 - ★ At temperatures closer to T_c , the contribution of heavier states becomes non-negligible
- ▶ While the lattice determination of heavy states in the spectrum is challenging [M. Teper, 1998], the spectral density $\hat{\rho}(m)$ can be predicted by modelling glueballs as “rings of glue” [N. Isgur and J. E. Paton, 1984]
- ▶ The effective description of confining flux tubes as open strings [M. Lüscher, 1980] suggests to model the glueballs as closed bosonic strings, leading to:

$$\hat{\rho}(m) = \frac{1}{m} \left(\frac{2\pi T_H}{3m} \right)^3 \exp(m/T_H)$$

where T_H is the Hagedorn temperature [R. Hagedorn, 1965]

- ▶ The total contribution to the pressure, including contributions from the lightest states and from a continuous spectrum above a threshold $m_{\text{threshold}} = 2m_{0^{++}}$, reads

$$p(T) = \sum_{m_i < m_{\text{threshold}}} (2J_i + 1) p(m_i; T) + n_C \int_{m_{\text{threshold}}}^{\infty} dm' \hat{\rho}(m') p(m'; T)$$

where n_C accounts for the charge-conjugation multiplicity



A string model for the spectrum

- ▶ How many glueball states should be included in the sum?
 - ★ At temperatures much lower than the deconfinement transition temperature T_c , the sum is completely dominated by the lightest glueballs
 - ★ At temperatures closer to T_c , the contribution of heavier states becomes non-negligible
- ▶ While the lattice determination of heavy states in the spectrum is challenging [M. Teper, 1998], the spectral density $\hat{\rho}(m)$ can be predicted by modelling glueballs as “rings of glue” [N. Isgur and J. E. Paton, 1984]
- ▶ The effective description of confining flux tubes as open strings [M. Lüscher, 1980] suggests to model the glueballs as closed bosonic strings, leading to:

$$\hat{\rho}(m) = \frac{1}{m} \left(\frac{2\pi T_H}{3m} \right)^3 \exp(m/T_H)$$

where T_H is the Hagedorn temperature [R. Hagedorn, 1965]

- ▶ The total contribution to the pressure, including contributions from the lightest states and from a continuous spectrum above a threshold $m_{\text{threshold}} = 2m_{0^{++}}$, reads

$$p(T) = \sum_{m_i < m_{\text{threshold}}} (2J_i + 1) p(m_i; T) + n_C \int_{m_{\text{threshold}}}^{\infty} dm' \hat{\rho}(m') p(m'; T)$$

where n_C accounts for the charge-conjugation multiplicity



SU(2) gauge theory

- ▶ SU(2) Yang-Mills theory is the simplest confining gauge theory based on a non-Abelian Lie group
- ▶ Due to its small dimension, it has been studied as a QCD prototype since the early days of lattice QCD [M. Creutz, 1980]
- ▶ Larger Lie groups have (several) SU(2) subgroups—a feature which is also used in simulation algorithms [N. Cabibbo and E. Marinari, 1982]
- ▶ The theory undergoes a *second-order* deconfining phase transition at a critical temperature T_c [J. Engels, F. Karsch and K. Redlich, 1994]
- ▶ When (a sufficiently large number of) dynamical fermions (in a suitable representation) are included, the theory could provide a potentially viable walking technicolor model for dynamical electro-weak symmetry breaking (DEWSB) [Del Debbio, 2010]
- ▶ When a small number of fundamental fermions are included, the (pseudo-)real nature of the SU(2) representations allows one to simulate the theory at finite quark chemical potential [S. Hands et al., 1999]
- ▶ The (pseudo-)real nature of the gauge group representations implies that in SU(2) Yang-Mills theory there exist no glueballs with $C = -1$ charge-conjugation eigenvalue



SU(2) gauge theory

- ▶ SU(2) Yang-Mills theory is the simplest confining gauge theory based on a non-Abelian Lie group
- ▶ Due to its small dimension, it has been studied as a QCD prototype since the early days of lattice QCD [M. Creutz, 1980]
- ▶ Larger Lie groups have (several) SU(2) subgroups—a feature which is also used in simulation algorithms [N. Cabibbo and E. Marinari, 1982]
- ▶ The theory undergoes a *second-order* deconfining phase transition at a critical temperature T_c [J. Engels, F. Karsch and K. Redlich, 1994]
- ▶ When (a sufficiently large number of) dynamical fermions (in a suitable representation) are included, the theory could provide a potentially viable walking technicolor model for dynamical electro-weak symmetry breaking (DEWSB) [Del Debbio, 2010]
- ▶ When a small number of fundamental fermions are included, the (pseudo-)real nature of the SU(2) representations allows one to simulate the theory at finite quark chemical potential [S. Hands et al., 1999]
- ▶ The (pseudo-)real nature of the gauge group representations implies that in SU(2) Yang-Mills theory there exist no glueballs with $C = -1$ charge-conjugation eigenvalue



SU(2) gauge theory

- ▶ SU(2) Yang-Mills theory is the simplest confining gauge theory based on a non-Abelian Lie group
- ▶ Due to its small dimension, it has been studied as a QCD prototype since the early days of lattice QCD [M. Creutz, 1980]
- ▶ Larger Lie groups have (several) SU(2) subgroups—a feature which is also used in simulation algorithms [N. Cabibbo and E. Marinari, 1982]
- ▶ The theory undergoes a *second-order* deconfining phase transition at a critical temperature T_c [J. Engels, F. Karsch and K. Redlich, 1994]
- ▶ When (a sufficiently large number of) dynamical fermions (in a suitable representation) are included, the theory could provide a potentially viable walking technicolor model for dynamical electro-weak symmetry breaking (DEWSB) [Del Debbio, 2010]
- ▶ When a small number of fundamental fermions are included, the (pseudo-)real nature of the SU(2) representations allows one to simulate the theory at finite quark chemical potential [S. Hands et al., 1999]
- ▶ The (pseudo-)real nature of the gauge group representations implies that in SU(2) Yang-Mills theory there exist no glueballs with $C = -1$ charge-conjugation eigenvalue



SU(2) gauge theory

- ▶ SU(2) Yang-Mills theory is the simplest confining gauge theory based on a non-Abelian Lie group
- ▶ Due to its small dimension, it has been studied as a QCD prototype since the early days of lattice QCD [M. Creutz, 1980]
- ▶ Larger Lie groups have (several) SU(2) subgroups—a feature which is also used in simulation algorithms [N. Cabibbo and E. Marinari, 1982]
- ▶ The theory undergoes a *second-order* deconfining phase transition at a critical temperature T_c [J. Engels, F. Karsch and K. Redlich, 1994]
- ▶ When (a sufficiently large number of) dynamical fermions (in a suitable representation) are included, the theory could provide a potentially viable walking technicolor model for dynamical electro-weak symmetry breaking (DEWSB) [Del Debbio, 2010]
- ▶ When a small number of fundamental fermions are included, the (pseudo-)real nature of the SU(2) representations allows one to simulate the theory at finite quark chemical potential [S. Hands et al., 1999]
- ▶ The (pseudo-)real nature of the gauge group representations implies that in SU(2) Yang-Mills theory there exist no glueballs with $C = -1$ charge-conjugation eigenvalue



SU(2) gauge theory

- ▶ SU(2) Yang-Mills theory is the simplest confining gauge theory based on a non-Abelian Lie group
- ▶ Due to its small dimension, it has been studied as a QCD prototype since the early days of lattice QCD [M. Creutz, 1980]
- ▶ Larger Lie groups have (several) SU(2) subgroups—a feature which is also used in simulation algorithms [N. Cabibbo and E. Marinari, 1982]
- ▶ The theory undergoes a *second-order* deconfining phase transition at a critical temperature T_c [J. Engels, F. Karsch and K. Redlich, 1994]
- ▶ When (a sufficiently large number of) dynamical fermions (in a suitable representation) are included, the theory could provide a potentially viable walking technicolor model for dynamical electro-weak symmetry breaking (DEWSB) [Del Debbio, 2010]
- ▶ When a small number of fundamental fermions are included, the (pseudo-)real nature of the SU(2) representations allows one to simulate the theory at finite quark chemical potential [S. Hands et al., 1999]
- ▶ The (pseudo-)real nature of the gauge group representations implies that in SU(2) Yang-Mills theory there exist no glueballs with $C = -1$ charge-conjugation eigenvalue



SU(2) gauge theory

- ▶ SU(2) Yang-Mills theory is the simplest confining gauge theory based on a non-Abelian Lie group
- ▶ Due to its small dimension, it has been studied as a QCD prototype since the early days of lattice QCD [M. Creutz, 1980]
- ▶ Larger Lie groups have (several) SU(2) subgroups—a feature which is also used in simulation algorithms [N. Cabibbo and E. Marinari, 1982]
- ▶ The theory undergoes a *second-order* deconfining phase transition at a critical temperature T_c [J. Engels, F. Karsch and K. Redlich, 1994]
- ▶ When (a sufficiently large number of) dynamical fermions (in a suitable representation) are included, the theory could provide a potentially viable walking technicolor model for dynamical electro-weak symmetry breaking (DEWSB) [Del Debbio, 2010]
- ▶ When a small number of fundamental fermions are included, the (pseudo-)real nature of the SU(2) representations allows one to simulate the theory at finite quark chemical potential [S. Hands et al., 1999]
- ▶ The (pseudo-)real nature of the gauge group representations implies that in SU(2) Yang-Mills theory there exist no glueballs with $C = -1$ charge-conjugation eigenvalue



SU(2) gauge theory

- ▶ SU(2) Yang-Mills theory is the simplest confining gauge theory based on a non-Abelian Lie group
- ▶ Due to its small dimension, it has been studied as a QCD prototype since the early days of lattice QCD [M. Creutz, 1980]
- ▶ Larger Lie groups have (several) SU(2) subgroups—a feature which is also used in simulation algorithms [N. Cabibbo and E. Marinari, 1982]
- ▶ The theory undergoes a *second-order* deconfining phase transition at a critical temperature T_c [J. Engels, F. Karsch and K. Redlich, 1994]
- ▶ When (a sufficiently large number of) dynamical fermions (in a suitable representation) are included, the theory could provide a potentially viable walking technicolor model for dynamical electro-weak symmetry breaking (DEWSB) [Del Debbio, 2010]
- ▶ When a small number of fundamental fermions are included, the (pseudo-)real nature of the SU(2) representations allows one to simulate the theory at finite quark chemical potential [S. Hands et al., 1999]
- ▶ The (pseudo-)real nature of the gauge group representations implies that in SU(2) Yang-Mills theory there exist no glueballs with $C = -1$ charge-conjugation eigenvalue

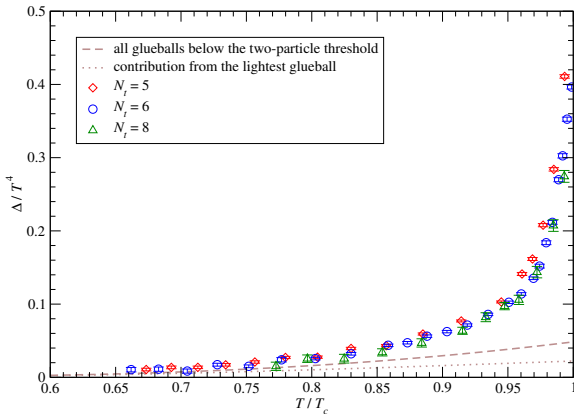


Lattice results: Impact of heavy states

Our results for the trace anomaly in units of the fourth power of the temperature

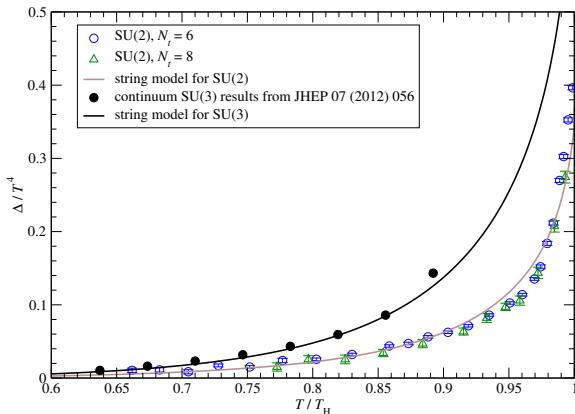
$$\frac{\Delta}{T^4} = T \frac{\partial}{\partial T} \left(\frac{p}{T^4} \right)$$

show that heavy glueballs give a large contribution to the equation of state close to T_c



Lattice results: Comparison with the bosonic string model

The predictions of a gas of massive glueballs, including contributions from heavy states modelled by a closed bosonic string model, are in excellent agreement with lattice data for both SU(2) and SU(3) Yang-Mills theories [S. Borsányi et al., 2012]—see also [H. B. Meyer, 2009], [M. Caselle et al., 2011], [F. Buisseret and G. Lacroix, 2011]



Outline

Introduction

Thermodynamics in the confining phase: comparing $SU(2)$ and $SU(3)$ Yang-Mills theories

Thermodynamics in the deconfined phase: comparing $SU(N)$ and G_2 theories

Conclusions



G₂ Yang-Mills theory

- ▶ The exceptional group G₂ (with dimension 14 and rank 2) is the smallest compact, simply connected Lie group with a trivial center
- ▶ “Quarks” in the fundamental representation **7**, “gluons” in the adjoint **14**
- ▶ Due to the absence of center vortices, which are important rôle for confinement [L. Del Debbio et al., 1996], [P. de Forcrand and M. D’Elia, 1999], G₂ Yang-Mills theory has triggered interest in the lattice community [K. Holland et al., 2003]
- ▶ The theory is actually “screening” at large distances: a fundamental color source can be screened by three gluons

$$\begin{aligned}
 14 \otimes 14 \otimes 14 = & \mathbf{1} \oplus \mathbf{7} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{64} \oplus \mathbf{64} \\
 & \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77}' \oplus \mathbf{77}' \oplus \mathbf{77}' \oplus \mathbf{182} \oplus \mathbf{189} \oplus \mathbf{189} \oplus \mathbf{189} \\
 & \oplus \mathbf{273} \oplus \mathbf{448} \oplus \mathbf{448}
 \end{aligned}$$

- ▶ The absence of *N*-ality also implies that the theory admits both bosonic and fermionic “baryons”

$$7 \otimes 7 = \mathbf{1} \oplus \mathbf{7} \oplus \mathbf{14} \oplus \mathbf{27}$$

- ▶ As an SO(7) subgroup, G₂ has real representations: this allows one to simulate the theory with dynamical quarks at finite chemical potential [A. Maas et al., 2012]
- ▶ The theory with dynamical fermions has also been studied as a candidate walking technicolor model for DEWSB [M. Mojaza et al., 2012]



G₂ Yang-Mills theory

- ▶ The exceptional group G₂ (with dimension 14 and rank 2) is the smallest compact, simply connected Lie group with a trivial center
- ▶ “Quarks” in the fundamental representation **7**, “gluons” in the adjoint **14**
- ▶ Due to the absence of center vortices, which are important rôle for confinement [L. Del Debbio et al., 1996], [P. de Forcrand and M. D’Elia, 1999], G₂ Yang-Mills theory has triggered interest in the lattice community [K. Holland et al., 2003]
- ▶ The theory is actually “screening” at large distances: a fundamental color source can be screened by three gluons

$$\begin{aligned}
 14 \otimes 14 \otimes 14 = & 1 \oplus 7 \oplus 14 \oplus 14 \oplus 14 \oplus 14 \oplus 14 \oplus 27 \oplus 27 \oplus 27 \oplus 64 \oplus 64 \\
 & \oplus 77 \oplus 77 \oplus 77 \oplus 77 \oplus 77' \oplus 77' \oplus 77' \oplus 182 \oplus 189 \oplus 189 \oplus 189 \\
 & \oplus 273 \oplus 448 \oplus 448
 \end{aligned}$$

- ▶ The absence of *N*-ality also implies that the theory admits both bosonic and fermionic “baryons”

$$7 \otimes 7 = 1 \oplus 7 \oplus 14 \oplus 27$$

- ▶ As an SO(7) subgroup, G₂ has real representations: this allows one to simulate the theory with dynamical quarks at finite chemical potential [A. Maas et al., 2012]
- ▶ The theory with dynamical fermions has also been studied as a candidate walking technicolor model for DEWSB [M. Mojaza et al., 2012]



G₂ Yang-Mills theory

- ▶ The exceptional group G₂ (with dimension 14 and rank 2) is the smallest compact, simply connected Lie group with a trivial center
- ▶ “Quarks” in the fundamental representation **7**, “gluons” in the adjoint **14**
- ▶ Due to the absence of center vortices, which are important rôle for confinement [L. Del Debbio et al., 1996], [P. de Forcrand and M. D’Elia, 1999], G₂ Yang-Mills theory has triggered interest in the lattice community [K. Holland et al., 2003]
- ▶ The theory is actually “screening” at large distances: a fundamental color source can be screened by three gluons

$$\begin{aligned}
 14 \otimes 14 \otimes 14 &= 1 \oplus 7 \oplus 14 \oplus 14 \oplus 14 \oplus 14 \oplus 14 \oplus 27 \oplus 27 \oplus 27 \oplus 64 \oplus 64 \\
 &\oplus 77 \oplus 77 \oplus 77 \oplus 77 \oplus 77' \oplus 77' \oplus 77' \oplus 182 \oplus 189 \oplus 189 \oplus 189 \\
 &\oplus 273 \oplus 448 \oplus 448
 \end{aligned}$$

- ▶ The absence of *N*-ality also implies that the theory admits both bosonic and fermionic “baryons”

$$7 \otimes 7 = 1 \oplus 7 \oplus 14 \oplus 27$$

- ▶ As an SO(7) subgroup, G₂ has real representations: this allows one to simulate the theory with dynamical quarks at finite chemical potential [A. Maas et al., 2012]
- ▶ The theory with dynamical fermions has also been studied as a candidate walking technicolor model for DEWSB [M. Mojaza et al., 2012]



G₂ Yang-Mills theory

- ▶ The exceptional group G₂ (with dimension 14 and rank 2) is the smallest compact, simply connected Lie group with a trivial center
- ▶ “Quarks” in the fundamental representation **7**, “gluons” in the adjoint **14**
- ▶ Due to the absence of center vortices, which are important rôle for confinement [L. Del Debbio et al., 1996], [P. de Forcrand and M. D’Elia, 1999], G₂ Yang-Mills theory has triggered interest in the lattice community [K. Holland et al., 2003]
- ▶ The theory is actually “screening” at large distances: a fundamental color source can be screened by three gluons

$$\begin{aligned}
 \mathbf{14} \otimes \mathbf{14} \otimes \mathbf{14} = & \mathbf{1} \oplus \mathbf{7} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{64} \oplus \mathbf{64} \\
 & \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77}' \oplus \mathbf{77}' \oplus \mathbf{77}' \oplus \mathbf{182} \oplus \mathbf{189} \oplus \mathbf{189} \oplus \mathbf{189} \\
 & \oplus \mathbf{273} \oplus \mathbf{448} \oplus \mathbf{448}
 \end{aligned}$$

- ▶ The absence of *N*-ality also implies that the theory admits both bosonic and fermionic “baryons”

$$\mathbf{7} \otimes \mathbf{7} = \mathbf{1} \oplus \mathbf{7} \oplus \mathbf{14} \oplus \mathbf{27}$$

- ▶ As an SO(7) subgroup, G₂ has real representations: this allows one to simulate the theory with dynamical quarks at finite chemical potential [A. Maas et al., 2012]
- ▶ The theory with dynamical fermions has also been studied as a candidate walking technicolor model for DEWSB [M. Mojaza et al., 2012]



G₂ Yang-Mills theory

- ▶ The exceptional group G₂ (with dimension 14 and rank 2) is the smallest compact, simply connected Lie group with a trivial center
- ▶ “Quarks” in the fundamental representation **7**, “gluons” in the adjoint **14**
- ▶ Due to the absence of center vortices, which are important rôle for confinement [L. Del Debbio et al., 1996], [P. de Forcrand and M. D’Elia, 1999], G₂ Yang-Mills theory has triggered interest in the lattice community [K. Holland et al., 2003]
- ▶ The theory is actually “screening” at large distances: a fundamental color source can be screened by three gluons

$$\begin{aligned}
 \mathbf{14} \otimes \mathbf{14} \otimes \mathbf{14} = & \mathbf{1} \oplus \mathbf{7} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{64} \oplus \mathbf{64} \\
 & \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77}' \oplus \mathbf{77}' \oplus \mathbf{77}' \oplus \mathbf{182} \oplus \mathbf{189} \oplus \mathbf{189} \oplus \mathbf{189} \\
 & \oplus \mathbf{273} \oplus \mathbf{448} \oplus \mathbf{448}
 \end{aligned}$$

- ▶ The absence of *N*-ality also implies that the theory admits both bosonic and fermionic “baryons”

$$\mathbf{7} \otimes \mathbf{7} = \mathbf{1} \oplus \mathbf{7} \oplus \mathbf{14} \oplus \mathbf{27}$$

- ▶ As an SO(7) subgroup, G₂ has real representations: this allows one to simulate the theory with dynamical quarks at finite chemical potential [A. Maas et al., 2012]
- ▶ The theory with dynamical fermions has also been studied as a candidate walking technicolor model for DEWSB [M. Mojaza et al., 2012]



G₂ Yang-Mills theory

- ▶ The exceptional group G₂ (with dimension 14 and rank 2) is the smallest compact, simply connected Lie group with a trivial center
- ▶ “Quarks” in the fundamental representation **7**, “gluons” in the adjoint **14**
- ▶ Due to the absence of center vortices, which are important rôle for confinement [L. Del Debbio et al., 1996], [P. de Forcrand and M. D’Elia, 1999], G₂ Yang-Mills theory has triggered interest in the lattice community [K. Holland et al., 2003]
- ▶ The theory is actually “screening” at large distances: a fundamental color source can be screened by three gluons

$$\begin{aligned}
 14 \otimes 14 \otimes 14 = & \mathbf{1} \oplus \mathbf{7} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{14} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{64} \oplus \mathbf{64} \\
 & \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77} \oplus \mathbf{77}' \oplus \mathbf{77}' \oplus \mathbf{77}' \oplus \mathbf{182} \oplus \mathbf{189} \oplus \mathbf{189} \oplus \mathbf{189} \\
 & \oplus \mathbf{273} \oplus \mathbf{448} \oplus \mathbf{448}
 \end{aligned}$$

- ▶ The absence of *N*-ality also implies that the theory admits both bosonic and fermionic “baryons”

$$7 \otimes 7 = \mathbf{1} \oplus \mathbf{7} \oplus \mathbf{14} \oplus \mathbf{27}$$

- ▶ As an SO(7) subgroup, G₂ has real representations: this allows one to simulate the theory with dynamical quarks at finite chemical potential [A. Maas et al., 2012]
- ▶ The theory with dynamical fermions has also been studied as a candidate walking technicolor model for DEWSB [M. Mojaza et al., 2012]



Lattice results: General features

- ★ First-order thermal deconfinement transition [G. Cossu et al., 2007], as predicted by semiclassical computations [E. Poppitz, T. Schäfer and M. Ünsal, 2012] and by phenomenological models [A. Dumitru et al., 2012]
- ★ The main equilibrium observables (pressure p , energy density ϵ and entropy density s) tend slowly towards their Stefan-Boltzmann limits



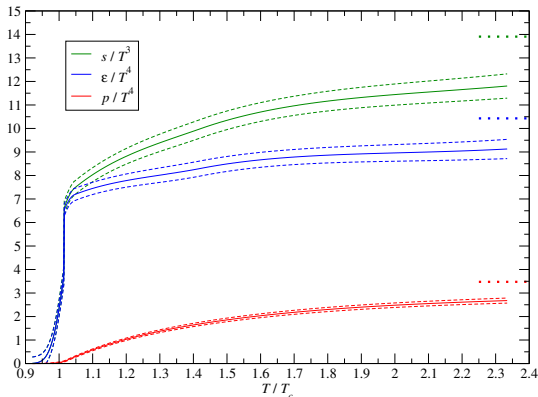
Lattice results: General features

- ★ First-order thermal deconfinement transition [G. Cossu et al., 2007], as predicted by semiclassical computations [E. Poppitz, T. Schäfer and M. Ünsal, 2012] and by phenomenological models [A. Dumitru et al., 2012]
- ★ The main equilibrium observables (pressure p , energy density ϵ and entropy density s) tend slowly towards their Stefan-Boltzmann limits



Lattice results: General features

- ★ First-order thermal deconfinement transition [G. Cossu et al., 2007], as predicted by semiclassical computations [E. Poppitz, T. Schäfer and M. Ünsal, 2012] and by phenomenological models [A. Dumitru et al., 2012]
- ★ The main equilibrium observables (pressure p , energy density ϵ and entropy density s) tend slowly towards their Stefan-Boltzmann limits



Lattice results: A closer look at quantitative details

- ★ In the deconfined phase, the thermodynamic observables *normalized per gluon d.o.f.* exhibit remarkable gauge-group independence—see also [M. P., 2009], [A. Mykkänen, M. P. and K. Rummukainen, 2012], [B. Lucini and M. P., 2012]
- ★ The trace anomaly Δ reveals a peculiar T^2 -dependence, probably of non-perturbative nature [R. D. Pisarski, 2006], also observed in 2 + 1 dimensions [M. Caselle et al., 2011]



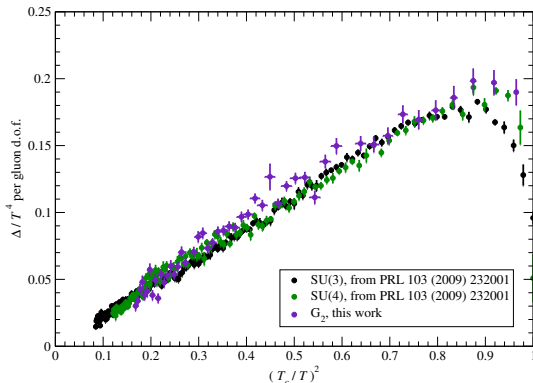
Lattice results: A closer look at quantitative details

- ★ In the deconfined phase, the thermodynamic observables *normalized per gluon d.o.f.* exhibit remarkable gauge-group independence—see also [M. P., 2009], [A. Mykkänen, M. P. and K. Rummukainen, 2012], [B. Lucini and M. P., 2012]
- ★ The trace anomaly Δ reveals a peculiar T^2 -dependence, probably of non-perturbative nature [R. D. Pisarski, 2006], also observed in 2 + 1 dimensions [M. Caselle et al., 2011]



Lattice results: A closer look at quantitative details

- ★ In the deconfined phase, the thermodynamic observables *normalized per gluon d.o.f.* exhibit remarkable gauge-group independence—see also [M. P., 2009], [A. Mykkänen, M. P. and K. Rummukainen, 2012], [B. Lucini and M. P., 2012]
- ★ The trace anomaly Δ reveals a peculiar T^2 -dependence, probably of non-perturbative nature [R. D. Pisarski, 2006], also observed in 2 + 1 dimensions [M. Caselle et al., 2011]



Outline

Introduction

Thermodynamics in the confining phase: comparing $SU(2)$ and $SU(3)$ Yang-Mills theories

Thermodynamics in the deconfined phase: comparing $SU(N)$ and G_2 theories

Conclusions



Concluding remarks

- ▶ Lattice QCD calculations are providing increasingly accurate first-principle predictions for QCD matter under extreme conditions
- ▶ Lattice simulations also reveal interesting universal, non-perturbative aspects in the thermodynamics of different, QCD-like gauge theories
 - ✧ The confining phase can be described as a gas of free glueballs, provided contributions from a Hagedorn-like excitation spectrum are included



Concluding remarks

- ▶ Lattice QCD calculations are providing increasingly accurate first-principle predictions for QCD matter under extreme conditions
- ▶ Lattice simulations also reveal interesting universal, non-perturbative aspects in the thermodynamics of different, QCD-like gauge theories
 - ★ The confining phase can be described as a gas of free glueballs, provided contributions from a Hagedorn-like excitation spectrum are included
 - ★ The EoS in the deconfined phase exhibits remarkable independence on the gauge group (up to a trivial color- and polarization-multiplicity factor) and a characteristic temperature dependence



Concluding remarks

- ▶ Lattice QCD calculations are providing increasingly accurate first-principle predictions for QCD matter under extreme conditions
- ▶ Lattice simulations also reveal interesting universal, non-perturbative aspects in the thermodynamics of different, QCD-like gauge theories
 - ★ The confining phase can be described as a gas of free glueballs, provided contributions from a Hagedorn-like excitation spectrum are included
 - ★ The EoS in the deconfined phase exhibits remarkable independence on the gauge group (up to a trivial color- and polarization-multiplicity factor) and a characteristic temperature dependence



Concluding remarks

- ▶ Lattice QCD calculations are providing increasingly accurate first-principle predictions for QCD matter under extreme conditions
- ▶ Lattice simulations also reveal interesting universal, non-perturbative aspects in the thermodynamics of different, QCD-like gauge theories
 - ★ The confining phase can be described as a gas of free glueballs, provided contributions from a Hagedorn-like excitation spectrum are included
 - ★ The EoS in the deconfined phase exhibits remarkable independence on the gauge group (up to a trivial color- and polarization-multiplicity factor) and a characteristic temperature dependence



Concluding remarks

- ▶ Lattice QCD calculations are providing increasingly accurate first-principle predictions for QCD matter under extreme conditions
- ▶ Lattice simulations also reveal interesting universal, non-perturbative aspects in the thermodynamics of different, QCD-like gauge theories
 - ★ The confining phase can be described as a gas of free glueballs, provided contributions from a Hagedorn-like excitation spectrum are included
 - ★ The EoS in the deconfined phase exhibits remarkable independence on the gauge group (up to a trivial color- and polarization-multiplicity factor) and a characteristic temperature dependence

Thanks for your attention!

